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# Power Laws in Natural (Critical) Phenomena

Mini-talk 4  
Examples of Power Laws

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# Agenda

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- Review of
  - Thermodynamic Functions and critical Indices
  - Scaling Hypothesis, Universality, Power laws
- Show and Tell
  - Physical
  - Other Science
- Next time:
  - Landau's Mean Field Theory

# Notable Persons of the Week

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# Origin of Thermodynamics

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$$Z = \sum_{\text{state}} e^{-\beta H} \equiv e^{\beta F} \quad \text{e.g. } H(\{s_i\}; J, H) = -J \sum \sum s_i s_j - \sum s_i H$$

*Derivatives of F*  $\Rightarrow$

- Equation of state:  $P = P(V, T, N)$
- Thermodynamic functions:  $C_H = -T \left. \frac{\partial^2 G(T, H)}{\partial T^2} \right|_H$
- Correlation:  $g(r, r') = \langle n(r) n(r') \rangle$
- Simplicity to Complexity
- Critical Phenomena, Complex?

# Parameters of Phase Transition

TABLE II. Parameters describing phase transition.

Physical quantity	Range of variables			Parameter describing quantity
	$\epsilon = (T - T_c)/T_c$	$h$	Behavior of quantity	
$\langle \phi \rangle$	$>0$	0	$\langle \phi \rangle = 0$	
	$<0$	0	$\langle \phi \rangle \sim \pm  \epsilon ^\beta$	$\beta$
	0	$\neq 0$	$\sim \pm  h ^{\beta/\alpha}$	$\beta/\alpha$
$\chi = \partial \langle \phi \rangle / \partial h _0$	$>0$	0	$\sim \epsilon^{-\gamma}$	$\gamma$
	$<0$	0	$\sim  \epsilon ^{-\gamma'}$	$\gamma'$
$g(r, r') = \langle \phi_r \phi_{r'} \rangle - \langle \phi \rangle^2$	0	0	$\sim  r - r' ^{-d+2-\eta}$	$\eta$
$\xi = \text{range of } g(r, r')$	$>0$	0	$\sim \epsilon^{-\nu}$	$\nu$
	$<0$	0	$\sim \epsilon^{-\nu'}$	$\nu'$
$C_h = \text{specific heat at constant } h$	$>0$	0	$\alpha \epsilon^{-\alpha} + b$	$\alpha$
	$<0$	0	$a'  \epsilon ^{-\alpha'} + b'$	$\alpha'$
	$\text{or}$	$>0$	$A \log \epsilon^{-1} + B$	$\alpha = 0$
		$<0$	$A' \log  \epsilon ^{-1} + B'$	$\alpha' = 0$

# Universality: an Illustration

Table XI. Correspondences between lattice gas model  
and Ising model.

Ising model	Lattice gas model
Number of up spins	Number of molecules
Partition function	Grand partition function
Number of spins	Volume
Free energy minus magnetic field	Pressure
Magnetization	Density minus critical density
Specific heat $C_m$	Specific heat $C_V$

# Critical Phenomena: Divergence of some thermodynamic Observables

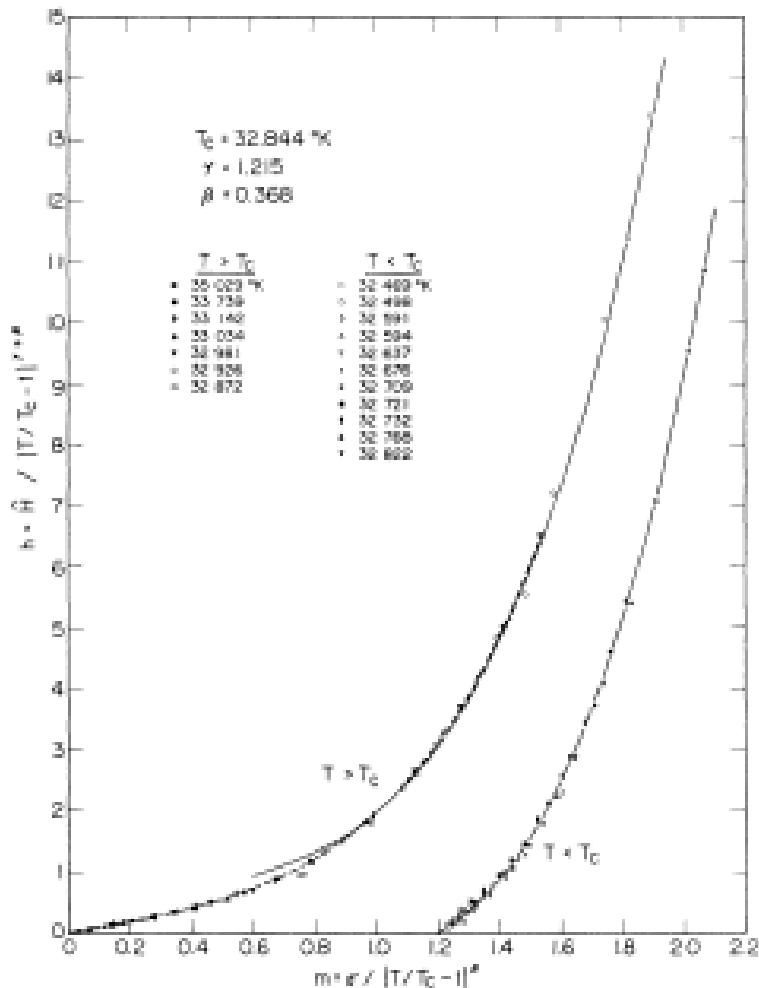


FIG. 1. A linear plot of scaled magnetic field  $h^*(\epsilon)$  as a function of scaled magnetization  $m^*$ . The dashed line represents Eq. (1), and fits experimental points from  $0 < m^* < 1.8$  when  $T > T_c$ . The solid lines are Eq. (2), and the dotted lines show Eq. (4) when it departs significantly from the experimental points. On this plot the origin is the critical isochore, the critical isotherm is at  $h = m^* = \infty$ , and the coexistence curve is represented by the point  $h = 0$ ,  $m^* \approx 1.2$ .

$$\frac{M}{\epsilon^\beta} = f\left(\frac{H}{\epsilon^{\beta\delta}}\right) ???$$

# Origin of Critical Indices

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- Widom's Scaling Hypothesis
- All thermodynamic functions (**Complexity**)
- Homogeneous functions of 2 variables
- Simplicity  $f(\lambda x, \lambda y) = \lambda^p f(x, y)$
- Relations among all indices



Figure 7: Michael Fisher, foreground, and Benjamin Widom, background, at a 65th birthday party for Michael.

# Critical Indices

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TABLE XIII. Definitions of critical indices for liquid-gas transition. If this transition is described correctly by the lattice gas model,  $\alpha$ ,  $\alpha'$ ,  $\beta$ ,  $\gamma$ ,  $\gamma'$ , and  $\delta$  all have the same values as in the Ising model.

$\epsilon$	$\epsilon = (T - T_c)/T_c$	
$\alpha'$	$C_v \sim (-\epsilon)^{-\alpha'}$	$\rho = \rho_{\text{L}}$ , $\epsilon < 0$
$\alpha$	$C_v \sim \epsilon^{-\alpha}$	$\rho = \rho_{\text{L}}$ , $\epsilon > 0$
$\beta$	$(\rho_L - \rho_G) \sim (-\epsilon)^{\beta}$	$\epsilon < 0$ , coexistence curve.
$\gamma'$	$K_T \sim (-\epsilon)^{-\gamma'}$	$\epsilon < 0$ , coexistence curve.
$\gamma$	$K_T \sim \epsilon^{-\gamma}$	$\rho = \rho_{\text{L}}$ , $\epsilon > 0$
$\delta$	$ P - P_c  \sim  \rho - \rho_c ^{\delta}$	$T = T_c$
$\eta$	Surface tension $\sim (-\epsilon)^{\eta}$	$\rho = \rho_{\text{L}}$ , $\epsilon < 0$

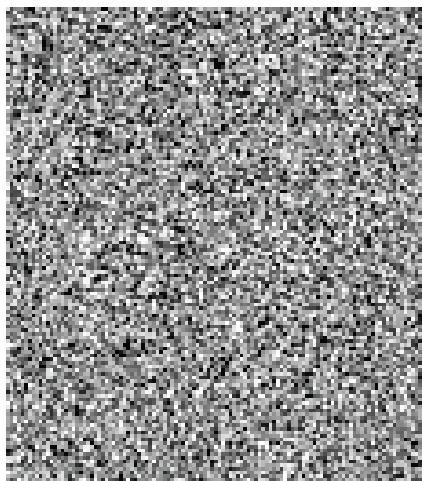
# Show and Tell

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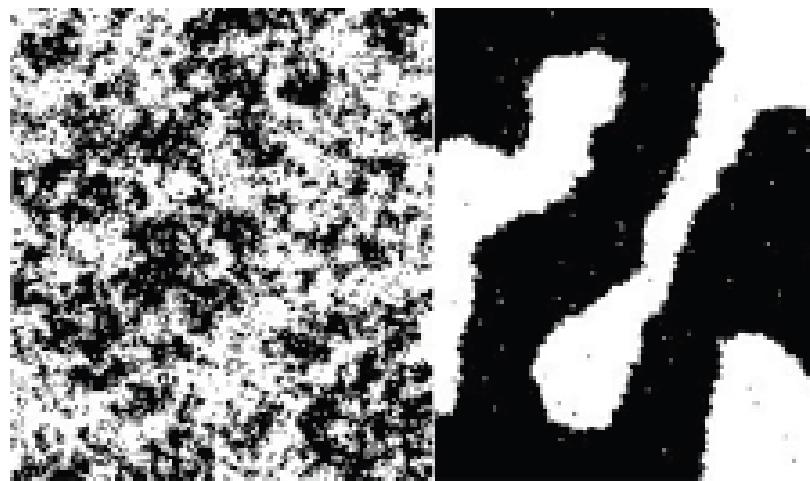
- More Pictures
- Less Equations (away from G....)

# 2D Ising model

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$B=0, J=0$

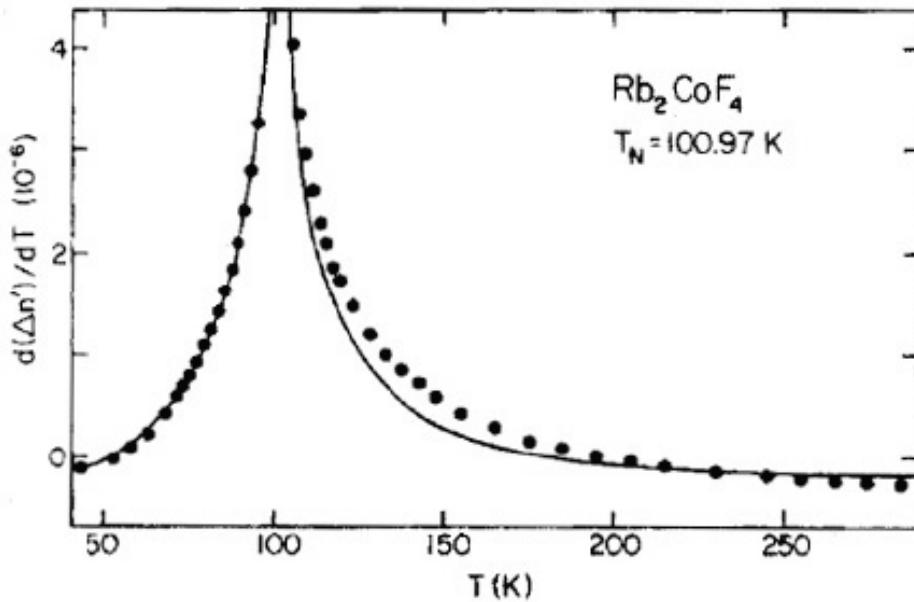


$B=0, J=0.4406868$

$B=0, J=0.8$

# Magnetic Specific Heat of $\text{Rb}_2\text{CoF}_4$

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$$c_n = -T \frac{\partial^2 G(T, H)}{\partial T^2} \Big|_H = A \frac{d(\Delta n)}{dT}$$

$n = \text{index of refraction}$

- $\text{Rb}_2\text{CoF}_4$  is a layered (2D Ising) antiferromagnet
- Nonmagnetic background ( $0.16\text{E}-6$ ) subtracted
- Amplitude adjusted for the unknown  $A$ .

# Superfluid $\lambda$ -Transition

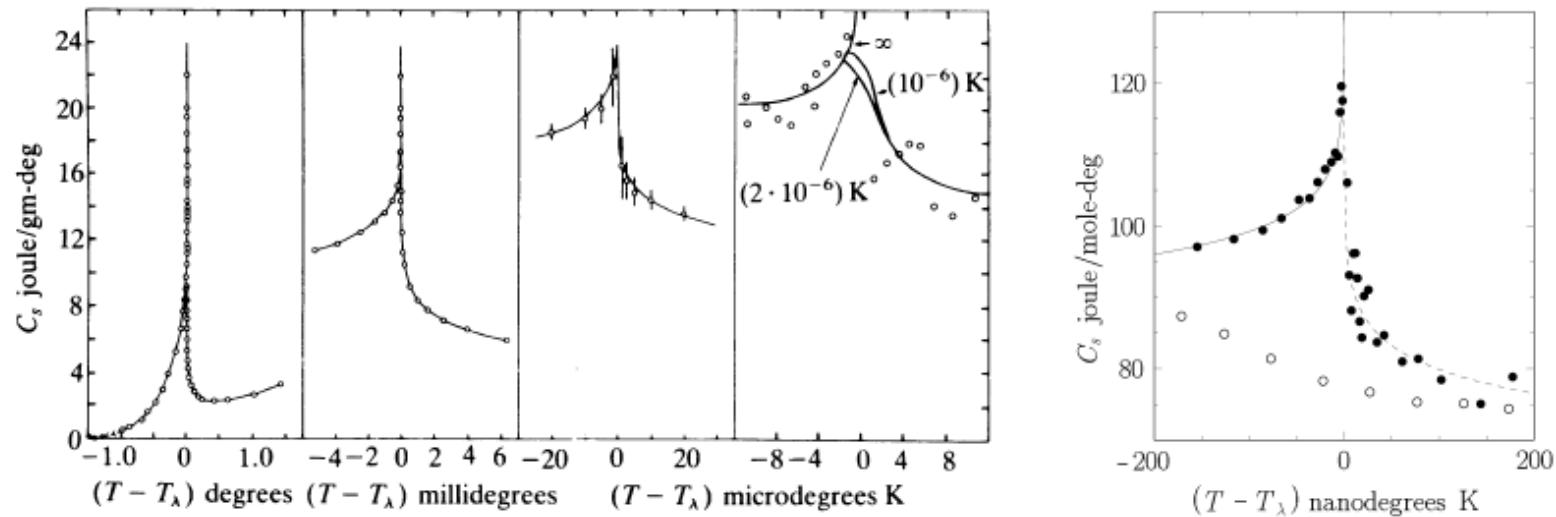


FIGURE 1.1 Specific heat near superfluid transition at  $T_\lambda \approx 2.18$  K measured with increasing temperature resolutions. The curve has the typical  $\lambda$ -shape which is the reason for calling it  $\lambda$ -transition. Note that at higher resolutions, the left shoulder of the peak lies above the right shoulder. The data are from Ref. [3]. The forth plot is broadened by the pressure difference between top and bottom of the sample. This is removed by the microgravity experiment in the space shuttle yielding the last plot (open circles are irrelevant here) [4]. They show no pressure broadening even in the nK regime around the critical temperature.

# $C_v$ of ${}^4\text{He}$ near $T_c$

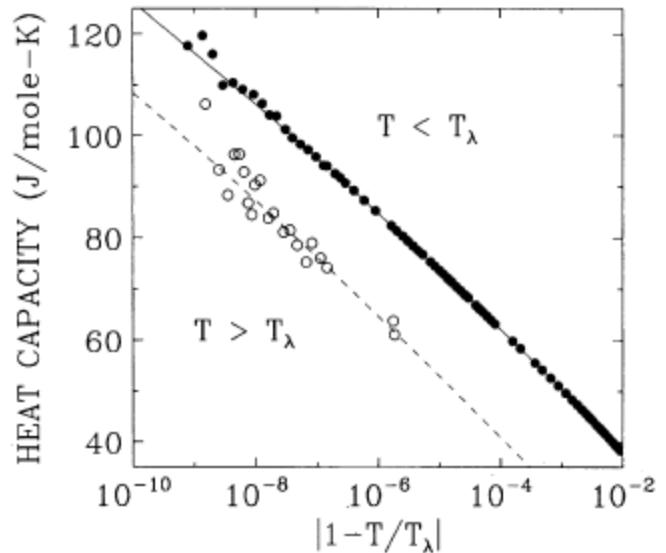
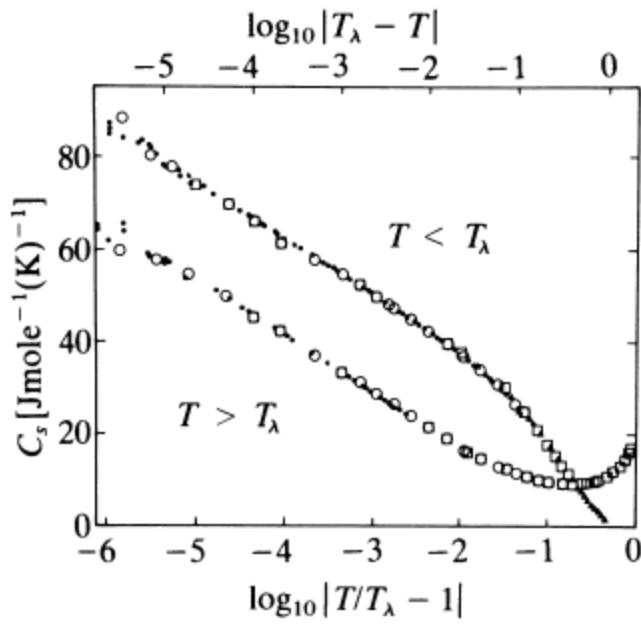


FIGURE 1.2 Specific heat of  ${}^4\text{He}$  near superfluid transition plotted against  $\log_{10}|T/T_\lambda - 1|$ . The early data on the left-hand side by G. Ahlers [5] yield  $\alpha \approx -0.026 \pm 0.004$  (for these an upper scale shows  $\log_{10}|T_\lambda - T|$  with  $T$  measured in units of K). The right-hand side shows the space shuttle data of J.A. Lipa et al. in Refs. [4, 6], which yield  $\alpha = -0.01056 \pm 0.0004$ .

Ahlers, PRA 3, 696 (1971); Lipa et al, PRL 76, 944 (1996)

# Superfluid Density near $T_c$

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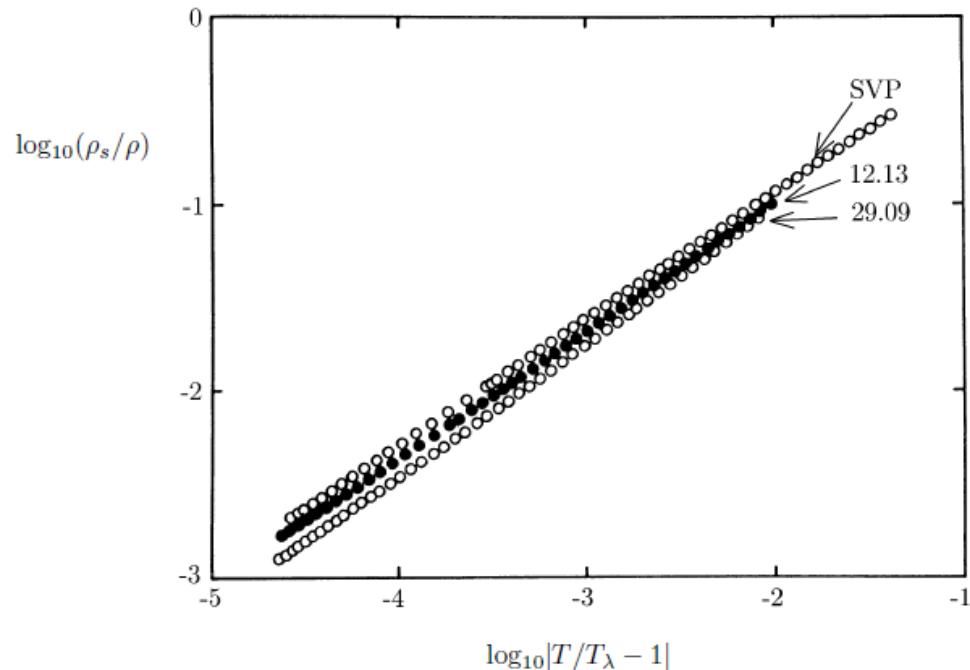


FIGURE 1.3 Doubly logarithmic plots of superfluid density  $\rho_s$  divided by the total density as a function of temperature. The slope is  $2/3$  [7]. The labels on the curves indicate the pressures in bar. The label SVP refers to the saturated vapor pressure.

Greywall and Ahlers, PRA 7, 2145 (1973)

# Fluid: Specific Heat of $\text{NiCl}_2 \cdot 6\text{H}_2\text{O}$

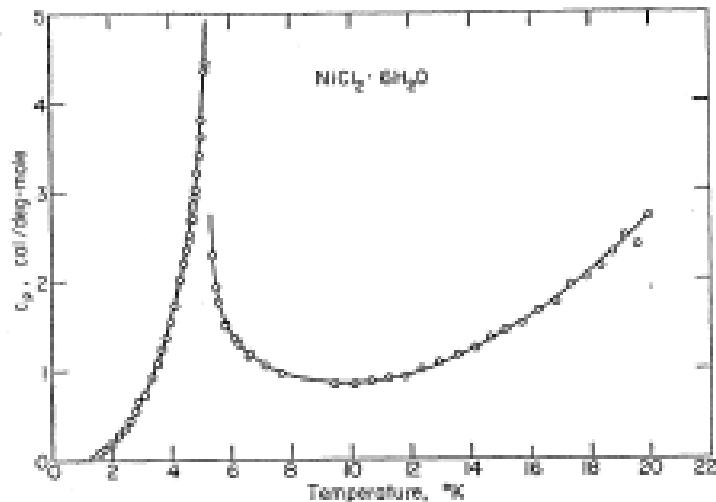


FIG. 1. Molar specific heat of  $\text{NiCl}_2 \cdot 6\text{H}_2\text{O}$  as a function of temperature.

- $t \equiv \frac{T - T_c}{T_c}$

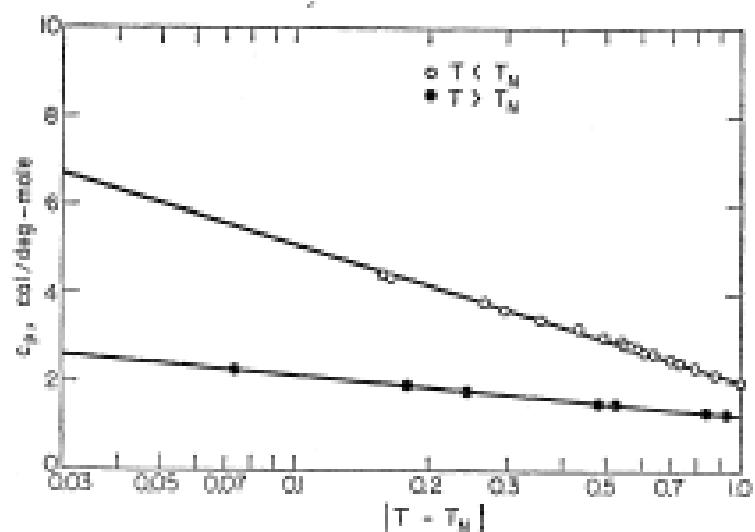


FIG. 5. Plot of  $C_{Hng}$  vs  $\log |\tau - \tau_N|$  for  $\text{NiCl}_2 \cdot 6\text{H}_2\text{O}$ .

- $C_p \sim t^{-\alpha} \quad t > 0$   
 $\sim (-t)^{-\alpha'} \quad t < 0$
- $\alpha < \alpha' ?$

# 1<sup>st</sup>: Relations among Indices?

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Relations among Critical Exponents

$$\alpha' + 2\beta + \gamma' \geq 2$$

$$\alpha' + \beta(\delta+1) \geq 2$$

$$(2-\alpha')\beta + 1 \geq (1-\alpha')\delta$$

$$\gamma'(\delta+1) \geq (2-\alpha')(\delta-1)$$

$$\gamma' \geq \beta(\delta-1)$$

:

$$(2-\alpha)\sigma \geq \delta_r + 1$$

!

$$d\nu \geq 2 - \alpha$$

Total of 17    (Griffiths - 1965b)  
(Stanley (1971))

# Scaling laws: predictions and Falsification

TABLE IV. Check of scaling law equalities for Ising model. The scaling laws predict that, for a given transition, all the numbers listed in the table (except for  $\Delta$  and  $y$ ) should be the same.

	$2-\alpha$	$2-\alpha'$	$d\tau$	$d\tau'$	$d\gamma/(2-\eta)$	$\gamma+2\beta$	$\gamma'+2\beta$	$\beta(\delta+1)$	$y$	$\Delta=2x/y$
Two-dimensional Ising model	2	2	2	2	2	2	2	2	1	3.75
Three-dimensional Ising model	1.87 $\pm 0.12$	1.93 + 0.04, - 0.16	1.93 $\pm 0.01$	?	1.933 $\pm 0.008$	1.87 $\pm 0.01$	1.94 $\pm 0.05$	1.93 $\pm 0.05$	1.55 $\pm 0.01$	3.22 $\pm 0.02$

TABLE X. Scaling law comparison.

	$2-\alpha$	$2-\alpha'$	$d\tau$	$d\tau'$	$d\gamma/2-\eta$	$\gamma+2\beta$	$\gamma'+2\beta$	$\beta(\delta+1)$	$y=1/\tau$	$\Delta^b$
Ferromagnets	1.92 $\pm 0.08$	1.92 $\pm 0.08$	1.95 $\pm 0.09$	...	2.08 $\pm 0.12$	1.99 $\pm 0.09$	...	1.7 $\pm 0.2$	1.54 $\pm 0.07$	3.2 $\pm 0.2$
Antiferromagnets	1.92 $\pm 0.08$	1.92 $\pm 0.08$	1.95 $\pm 0.09$	...	...	1.96 $\pm 0.10^a$	...	...	1.54 $\pm 0.07$	3.2 $\pm 0.2$

Kadanoff et al, RMP 39, 395 (1967)

# Fluid Compressibility

TABLE XVI. Value of  $\gamma$ .

Fluid	Value of $\gamma$	Reference	Range of $\epsilon$	Comments
Xe	$1.3 \pm 0.2$	99	$3 \times 10^{-4} < \epsilon < 3 \times 10^{-3}$	$\epsilon \leq 10^{-4}$ neglected
Ar	$0.6 \pm 0.2$	110	$3 \times 10^{-4} < \epsilon < 6 \times 10^{-3}$	Is extrapolation procedure right?
CO <sub>2</sub>	$1.37 \pm 0.2$	105, 104, 111	$10^{-3} < \epsilon < 10^{-2}$	
CO <sub>2</sub>	$1.0 \pm 0.1$	106	$10^{-4} < \epsilon < 2 \times 10^{-3}$	Perhaps pressure gradients are important
"Best value"	$1.37 \pm 0.2$			
Lattice gas	$1.250 \pm 0.001$	Table III		

TABLE XVII. Values of  $\gamma'$ .

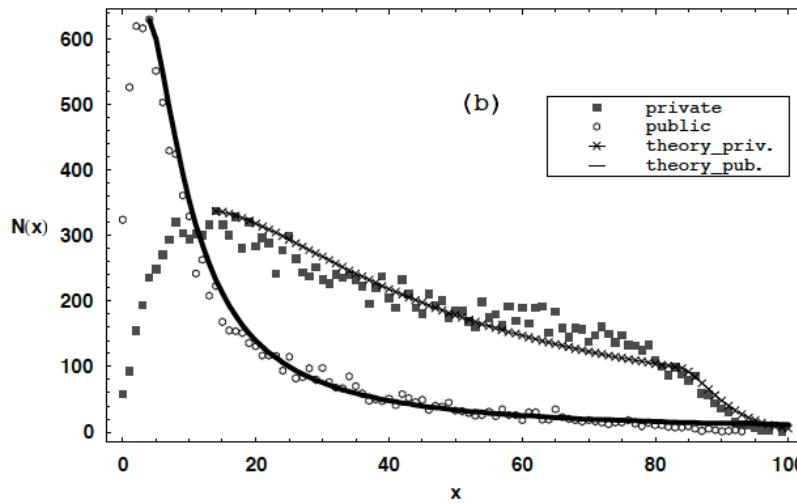
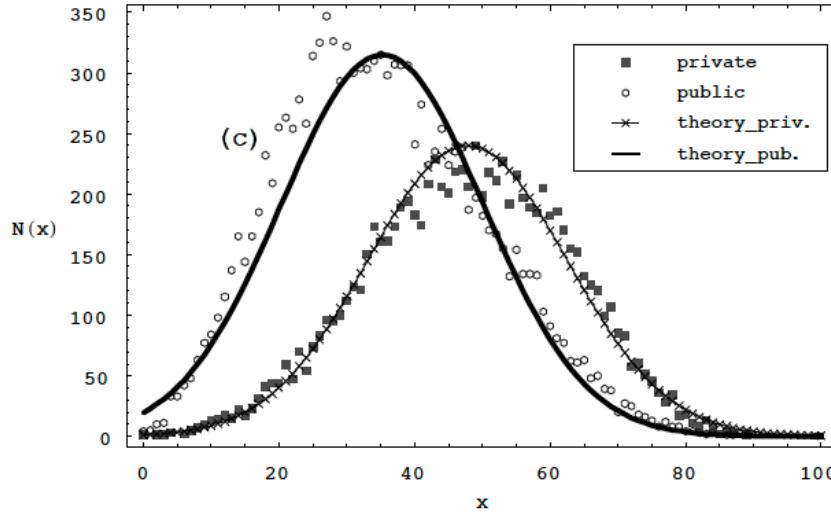
Fluid	Value of $\gamma'$	Reference	Range of $\epsilon$	Comments
Ar	$1.1 \pm 0.2$	110	$6 \times 10^{-4} < -\epsilon < 6 \times 10^{-3}$	Is extrapolation procedure right?
CO <sub>2</sub>	$1.0 \pm 0.3$	105, 104, 111	$\epsilon \approx 3 \times 10^{-3}$ $\epsilon \approx 3 \times 10^{-2}$ $\epsilon \approx 3 \times 10^{-4}$	If pressure gradients are important, this analysis gives no information.
"Best value"	$1.0 \pm 0.3$			Very little information
Lattice gas	$1.31 \pm 0.05$	Table III		

# Even

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- More Pictures
- Less equations

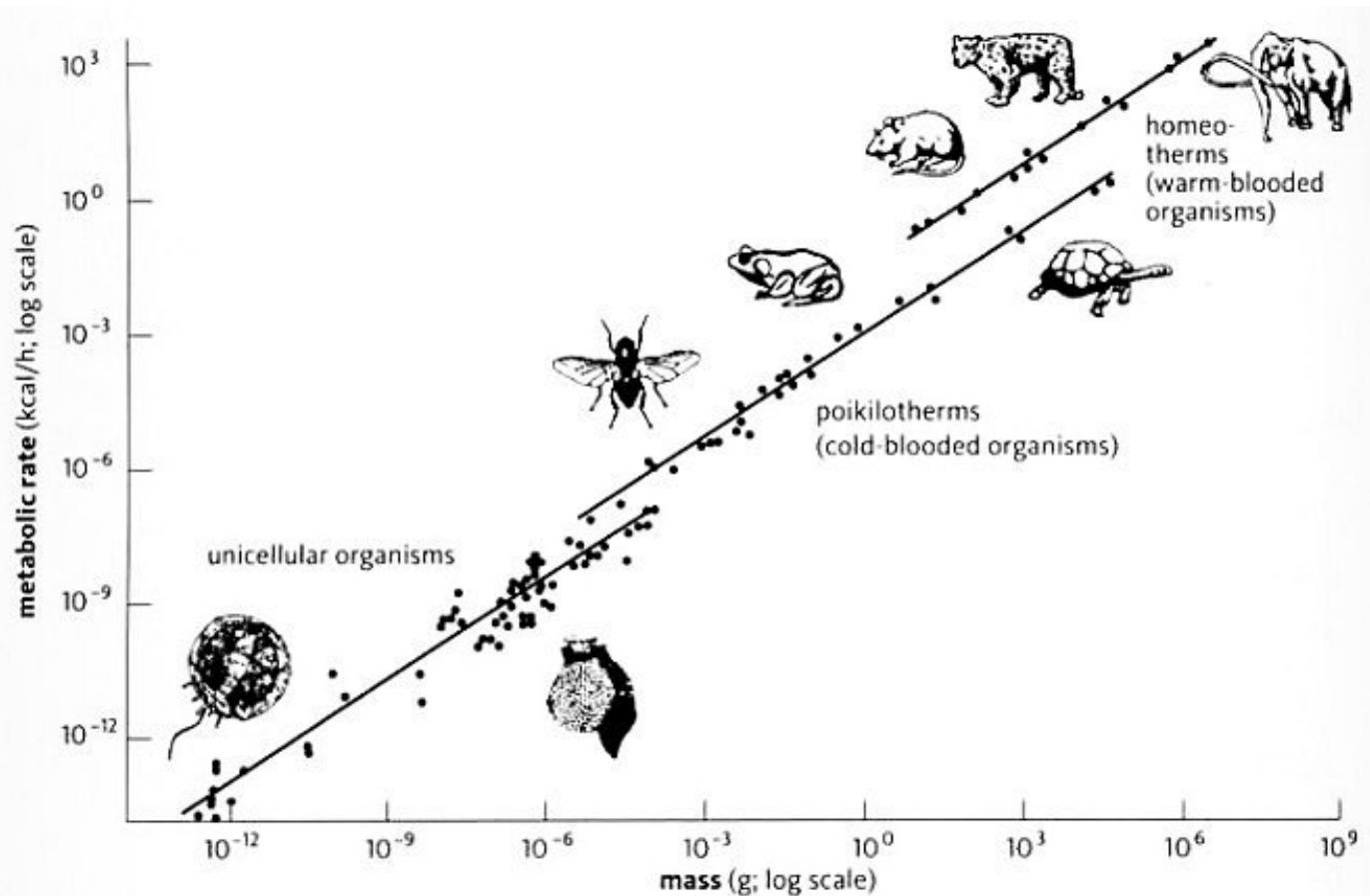
# Scaling in Sociology



$$p(x; \mu, \sigma^2) = \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{1}{2}\left(\frac{x-\mu}{\sigma}\right)^2}$$

$$p(x; \beta, T) = \frac{(\beta-1)T^{\beta-1}}{(x+T)^\beta}$$

# Scaling in Biology

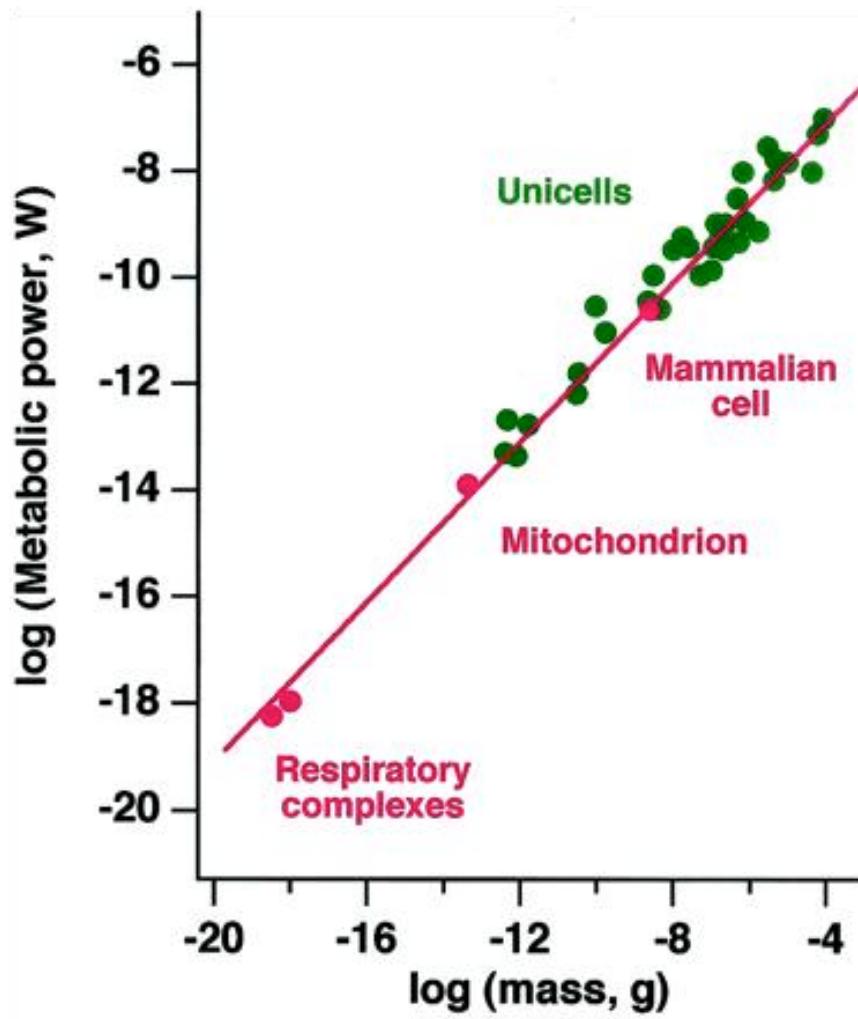


1 kcal/h = 1.162 watts

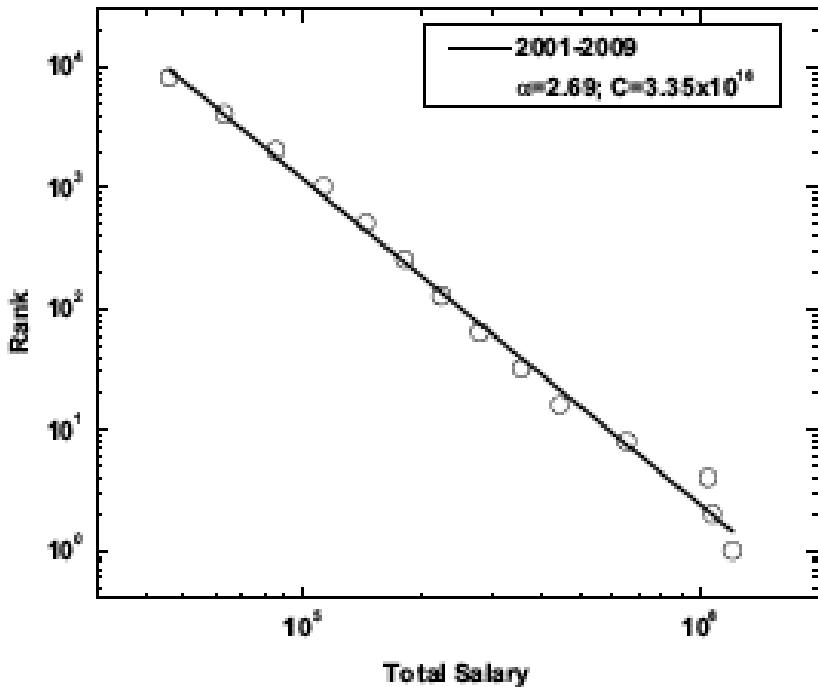
$$k \sim m^\alpha \quad (\alpha = 3/4)$$

Wikipedia

# Scaling in Microbiology



# Pareto: Income



$$cdf(x) = cx^{-\alpha}$$

$$p(x) \sim \frac{1}{x^{1+\alpha}}$$

Italian [economist Vilfredo Pareto](#), is a [power law probability distribution](#) that coincides with [social](#), [scientific](#), [geophysical](#), [actuarial](#), and many other types of observable phenomena.

# End Show and Tell

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- More Pictures
- Less equations